# Equivalence of Light-Front and Covariant Approaches in Meson-Baryon Interactions

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## Motivation

- Guidance from Approximate Chiral Symmetry of QCD -Extrapolation of Lattice Results for Small Pion Mass -Pion Cloud Contribution in Chiral Perturbation Theory (X<sub>PT</sub>) -Leading Non-Analytic (LNA) Behavior of Goldstone Boson Loops: Gell-Mann – Oaks – Renner Relation:  $m_{\pi}^2 f_{\pi}^2 = 2m_q < \overline{q}q >$
- Great Interest in Generalized Parton Distributions (GPDs)
   -Most Natural in Light-Front Dynamics (LFD)
  - -Infinite Momentum vs. Rest Frame as well as
    Pseudovector(PV) vs. Pseudoscalar(PS) Coupling in
    Equal-Time Calculations and Subtle Discrepancy between Results
    A.Thomas, W.Melnitchouk, F.Steffens, PRL85, 2892(00);
    J.-W. Chen, X.Ji, PRL87, 152002(01);
    M.Burkardt, private communication, etc.
- Discussion on Chiral Symmetry in LFD
  - -Distinguished Features: Vacuum, Symmetry
  - -Care of Treacherousness: Zero-Modes, End-Point Singularities, ...

# Outline

- Distinguished Features in LFD
  - Energy-Momentum Dispersion Relation
  - Common Belief of Equivalence
  - Treacherous Amplitudes
- Nucleon Self-Energy in  $\chi_{PT}$  (LNA Behaviors)
  - Manifestly Covariant Approach
  - Equal-Time Approach
    - Rest Frame vs. Infinite Momentum Frame
  - LFD Approach
- Conclusion and Outlook

#### **Distinguished Features in LFD**



**Energy-Momentum Dispersion Relations** 





**Manifestly Covariant Formulation** 



However, the proof of equivalence is treacherous.

Example: Arc-contribution in physical form factor B.Bakker, M.DeWitt, C.Ji, Y.Mishchenko, PRD72, 076005(05)



With the arc contribution, we confirmed the equivalence between covariant and light-front approach.

Treacherous Amplitudes in Nucleon Self-Energy C.Ji, W.Melnitchouk, A.Thomas, PRD80, 054018(09)

$$I = \int d^2k \frac{1}{k^2 - m^2 + i\epsilon}$$

Manifestly Covariant Calculation

$$I = -i\pi \left(\frac{2-n}{2} - \log \pi - \gamma + O(2-n) - \log \frac{m^2}{\mu^2}\right)_{n \to 2} \qquad I_{LNA} = i\pi \log m^2$$

**Time-Ordered Perturbation Theory** 

$$I = \int dk^{3} dk^{0} \frac{1}{(k^{0} + \omega_{k} - i\varepsilon)(k^{0} - \omega_{k} + i\varepsilon)}$$
  
=  $-2i\pi \int_{0}^{\Lambda \to \infty} dk^{3} \frac{1}{\sqrt{(k^{3})^{2} + m^{2}}}$   
=  $-2i\pi \log \left(\frac{2\Lambda}{m}\right)_{\Lambda \to \infty}$   
$$\omega_{k} - i\varepsilon = \sqrt{(k^{3})^{2} + m^{2}} - i\varepsilon$$



LF Polar Coordinate  $k^+ = r \cos \phi$   $k^- = r \sin \phi$ 

$$I = \int_{0}^{\infty} dr \, r \int_{0}^{2\pi} d\phi \frac{1}{r^{2} \sin \phi \cos \phi - m^{2} + i\varepsilon} = -i\pi \left( \log \frac{R^{2} e^{-i\pi/2}}{2m^{2}} + O(1/R^{4}) \right)_{R \to \infty}$$
$$I_{LNA}^{LF} = i\pi \log m^{2}$$

### $\pi N$ Interaction with PS vs. PV Coupling

$$L_{PS} = -g_{\pi NN} (\overline{\psi}_{N} i \gamma_{5} \vec{\tau} \psi_{N}) \cdot \vec{\phi}_{\pi} ; \quad L_{PV} = \frac{f_{\pi NN}}{m_{\pi}} (\overline{\psi}_{N} \gamma^{\mu} \gamma_{5} \vec{\tau} \psi_{N}) \cdot \partial_{\mu} \vec{\phi}_{\pi}$$
$$\frac{g_{\pi NN}}{2M} = \frac{f_{\pi NN}}{m_{\pi}} , \quad \frac{g_{A}}{f_{\pi}} = \frac{g_{\pi NN}}{M} \quad \text{(Goldberger-Treiman Relation)}$$

Coupling Constants;  $g_A \approx 1.267$ ,  $f_\pi \approx 93 \, MeV$ ,  $\frac{g_{\pi NN}^2}{4\pi} \approx 14.3$ ,  $\frac{f_{\pi NN}^2}{4\pi} \approx 0.08$ 

PV PV 
$$\Sigma^{PV} = \frac{1}{2} \sum_{s} \overline{u}(p,s) \hat{\Sigma}^{PV} u(p,s)$$

$$\hat{\Sigma}^{PV} = -i \left(\frac{2g_A}{f_\pi}\right)^2 \vec{\tau} \cdot \vec{\tau} \int \frac{d^4k}{(2\pi)^4} \frac{k\gamma_5(p-k+M)\gamma_5k}{D_\pi D_N}$$

 $D_{\pi} = k^2 - m_{\pi}^2 + i\varepsilon \qquad D_N = (p-k)^2 - M^2 + i\varepsilon$ 

### πN Interaction with PS vs. PV Coupling

$$L_{PS} = -g_{\pi NN} (\overline{\psi}_N i \gamma_5 \vec{\tau} \psi_N) \cdot \vec{\phi}_{\pi} \quad ; \quad L_{PV} = \frac{f_{\pi NN}}{m_{\pi}} (\overline{\psi}_N \gamma^{\mu} \gamma_5 \vec{\tau} \psi_N) \cdot \partial_{\mu} \vec{\phi}_{\pi}$$
$$g_{\pi NN} \quad f_{\pi NN} \quad g_A \quad g_{\pi NN}$$

$$\frac{\partial MN}{\partial M} = \frac{\partial MN}{m_{\pi}}, \quad \frac{\partial A}{f_{\pi}} = \frac{\partial MN}{M}$$
 (Goldberger-Treiman Relation)

Coupling Constants;  $g_A \approx 1.267$ ,  $f_\pi \approx 93 \, MeV$ ,  $\frac{g_{\pi NN}^2}{4\pi} \approx 14.3$ ,  $\frac{f_{\pi NN}^2}{4\pi} \approx 0.08$ 



$$\hat{\Sigma}^{PS} = -i \left(\frac{g_A M}{f_\pi}\right)^2 \vec{\tau} \cdot \vec{\tau} \int \frac{d^4 k}{(2\pi)^4} \frac{\gamma_5 (\not p - \not k + M) \gamma_5}{D_\pi D_N}$$

 $D_{\pi} = k^2 - m_{\pi}^2 + i\varepsilon \qquad D_N = (p-k)^2 - M^2 + i\varepsilon$ 

### **Nucleon Self-Energy**

$$\hat{\Sigma} = \Sigma_V p + \Sigma_S \implies \Sigma = M \Sigma_V + \Sigma_S$$

Substitution Technique:  $k^2 \rightarrow D_{\pi} + m_{\pi}^2$ ;  $p \cdot k \rightarrow \frac{1}{2}(D_{\pi} - D_N + m_{\pi}^2)$ 



$$M\Sigma_{V}^{PV} = \Sigma_{S}^{PV} = \frac{1}{2}\Sigma^{PV} = -\frac{3ig_{A}^{2}}{4f_{\pi}^{2}}\int \frac{d^{4}k}{(2\pi)^{4}}M\left[\frac{1}{D_{N}} + \frac{m_{\pi}^{2}}{D_{\pi}D_{N}}\right]$$

# Nucleon Self-Energy $\hat{\Sigma} = \Sigma_V p + \Sigma_S \implies \Sigma = M \Sigma_V + \Sigma_S$

Substitution Technique:  $k^2 \rightarrow D_{\pi} + m_{\pi}^2$ ;  $p \cdot k \rightarrow \frac{1}{2}(D_{\pi} - D_N + m_{\pi}^2)$ 



$$\begin{split} \Sigma^{PS} &= -\frac{3ig_A^2 M}{2f_\pi^2} \int \frac{d^4 k}{(2\pi)^4} \left[ \frac{1}{D_N} - \frac{1}{D_\pi} + \frac{m_\pi^2}{D_\pi D_N} \right]; \\ \Sigma^{PS}_V &= -3i \left( \frac{Mg_A}{f_\pi} \right)^2 \int \frac{d^4 k}{(2\pi)^4} \frac{1}{D_\pi D_N} , \\ \Sigma^{PS}_S &= -3i \left( \frac{Mg_A}{f_\pi} \right)^2 \int \frac{d^4 k}{(2\pi)^4} \frac{1}{2M} \left[ \frac{1}{D_N} - \frac{1}{D_\pi} + \frac{m_\pi^2 - 2M^2}{D_\pi D_N} \right] \end{split}$$

### **Manifestly Covariant Calculation**



$$\int d^4k \frac{1}{D_{\pi}D_N} = -i\pi^2 \left[ \gamma + \log \pi - \frac{2}{4-n} + \int_0^1 dx \log \frac{(1-x)^2 M^2 + x m_{\pi}^2}{\mu^2} + O(4-n) \right]_{n \to 4}$$

$$\int_{0}^{1} dx \log \frac{(1-x)^{2} M^{2} + x m_{\pi}^{2}}{\mu^{2}} = \log \frac{M^{2}}{\mu^{2}} + \frac{m_{\pi}^{2}}{2M^{2}} \log \frac{m_{\pi}^{2}}{M^{2}} + \frac{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}}{M^{2}} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}^{2} - 2M^{2}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{\sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} - \tan^{-1} \frac{m_{\pi}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^{2}}} \right) - 2m_{\pi}^{2} \left( \tan^{-1} \frac{m_{\pi}}{m_{\pi} \sqrt{4M^{2} - m_{\pi}^$$

LNA Behavior as  $m_{\pi} \rightarrow 0$ 

$$\tan^{-1} \frac{m_{\pi}^2 - 2M^2}{m_{\pi}\sqrt{4M^2 - m_{\pi}^2}} = -\frac{\pi}{2} + m_{\pi} + \frac{m_{\pi}^3}{24M^3} + O(\frac{m_{\pi}^5}{M^5}),$$
$$\frac{m_{\pi}\sqrt{4M^2 - m_{\pi}^2}}{M^2} = \frac{2m_{\pi}}{M} \left(1 - \frac{m_{\pi}^2}{8M^2} + O(\frac{m_{\pi}^4}{M^4})\right), \cdots$$

$$\Sigma^{PV} = -\frac{3ig_A^2}{2f_\pi^2} \int \frac{d^4k}{(2\pi)^4} M \left[\frac{1}{D_N} + \frac{m_\pi^2}{D_\pi D_N}\right]$$

$$\Sigma_{LNA}^{PV} = -\frac{3g_A^2}{32\pi f_\pi^2} \left[ m_\pi^3 + \frac{m_\pi^4}{2\pi M} \log m_\pi^2 + O(m_\pi^5) \right]$$

Contribution from a  $\Delta$  intermediate state  $\propto$ 

$$\frac{m_{\pi}^4}{M_{\Delta} - M} \log m_{\pi}^2$$

D.Leinweber, A.Thomas, K.Tsushima, S.Wright, PRD61,074502(00)

LNA Behavior as  $m_{\pi} \rightarrow 0$ 

$$\tan^{-1} \frac{m_{\pi}^2 - 2M^2}{m_{\pi}\sqrt{4M^2 - m_{\pi}^2}} = -\frac{\pi}{2} + m_{\pi} + \frac{m_{\pi}^3}{24M^3} + O(\frac{m_{\pi}^5}{M^5}),$$
$$\frac{m_{\pi}\sqrt{4M^2 - m_{\pi}^2}}{M^2} = \frac{2m_{\pi}}{M} \left(1 - \frac{m_{\pi}^2}{8M^2} + O(\frac{m_{\pi}^4}{M^4})\right), \cdots$$

$$\Sigma^{PS} = -\frac{3ig_A^2 M}{2f_\pi^2} \int \frac{d^4 k}{(2\pi)^4} \left[ \frac{1}{D_N} - \frac{1}{D_\pi} + \frac{m_\pi^2}{D_\pi D_N} \right]$$
$$\Sigma^{PS}_{LNA} = \frac{3g_A^2}{32\pi^2 f_\pi^2} \left[ \frac{M}{\pi} m_\pi^2 \log m_\pi^2 - m_\pi^3 + \frac{m_\pi^4}{2M} \log m_\pi^2 + O(m_\pi^5) \right]$$

$$M\Sigma_{V,LNA}^{PS} = \frac{3g_A^2 M}{32\pi^2 f_\pi^2} \left[ 2M^2 m_\pi + \frac{2M}{\pi} m_\pi^2 \log m_\pi^2 - m_\pi^3 + \cdots \right]$$
$$\Sigma_{S,LNA}^{PS} = -\frac{3g_A^2}{32\pi^2 f_\pi^2} \left[ 2M^2 m_\pi + \frac{M}{\pi} m_\pi^2 \log m_\pi^2 + \cdots \right]$$

## **Time-Ordered Perturbation Theory**



 $\Sigma_{ET}^{(++)} \quad \frac{\omega_k - i\varepsilon}{X} \quad E + E' - i\varepsilon}{X} \quad \Sigma_{ET}^{(--)} \frac{X}{-\omega_k + i\varepsilon} \quad \frac{X}{E - E' + i\varepsilon}$ 

### **Equal-Time Results**

$$\Sigma_{ET}^{(+-)LNA} = -\frac{3g_A^2}{32\pi f_\pi^2} \left[ m_\pi^3 + \frac{3m_\pi^4}{4\pi M} \log m_\pi^2 + O(m_\pi^5) \right]$$
$$\Sigma_{ET}^{(-+)LNA} = -\frac{3g_A^2}{32\pi f_\pi^2} \left[ -\frac{m_\pi^4}{4\pi M} \log m_\pi^2 + O(m_\pi^5) \right]$$

Without using the substitution technique, we get

$$\begin{split} \Sigma_{ET}^{\prime(+-)LNA} &= -\frac{3g_A^2}{32\pi f_\pi^2} \bigg[ m_\pi^3 + \frac{13m_\pi^4}{16\pi M} \log m_\pi^2 + O(m_\pi^5) \bigg] \\ \Sigma_{ET}^{\prime(-+)LNA} &= -\frac{3g_A^2}{32\pi f_\pi^2} \bigg[ -\frac{1}{2\pi} M m_\pi^2 \log m_\pi^2 - \frac{5m_\pi^4}{16\pi M} \log m_\pi^2 + O(m_\pi^5) \bigg] \\ \Sigma_{ET}^{\prime(++)LNA} &= -\frac{3g_A^2}{32\pi f_\pi^2} \bigg[ \frac{1}{2\pi} M m_\pi^2 \log m_\pi^2 + \frac{m_\pi^4}{8\pi M} \log m_\pi^2 + O(m_\pi^5) \bigg] \\ \Sigma_{ET}^{\prime(--)LNA} &= -\frac{3g_A^2}{32\pi f_\pi^2} \bigg[ \frac{1}{2\pi} M m_\pi^2 \log m_\pi^2 + \frac{m_\pi^4}{8\pi M} \log m_\pi^2 + O(m_\pi^5) \bigg] \end{split}$$

### Infinite Momentum Frame



## **Light-Front Dynamics**



## Conclusion and Outlook

 Equivalence of the covariant, equal-time (in both the rest frame and in the IMF), and LF formalisms is highly nontrivial.

-Arc contributions, moving poles, and end-point singularities.

 LNA behaviors of the vertex renormalization and the moments of the twist-two parton distribution functions are underway.

-Form factors, PDFs and GPDs.

• Hope to pave the way for a consistent interpretation of the physics of the pion cloud at the partonic level.